

# Ab Initio Prediction of the Structure and Vibration–Rotation Spectroscopic Properties of $\text{Li}_2\text{OH}$ and $\text{Li}_2\text{OH}^+$

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The equilibrium structures and potential energy surfaces of the  $\text{Li}_2\text{OH}$  radical and the  $\text{Li}_2\text{OH}^+$  cation in their ground electronic states have been determined from accurate ab initio calculations. The vibration–rotation energy levels and spectroscopic constants of three isotopic species ( $\text{Li}_2\text{OH}$ ,  $\text{Li}_2\text{OD}$ ,  $^6\text{Li}_2\text{OH}$ ) were calculated by a perturbational approach. The predicted spectroscopic constants may serve as a useful guide for detecting these species by vibration–rotation spectroscopy and for assigning their spectra.

## I. Introduction

Dilithium hydroxide,  $\text{Li}_2\text{OH}$ , is one of simplest examples of hypermetalated molecules.<sup>1</sup> The  $\text{Li}_2\text{OH}$  radical can be considered to result from reaction of a  $\text{LiOH}$  molecule with a Li atom. The extra Li atom is expected to bind to both O and Li atoms of  $\text{LiOH}$ . The isolated  $\text{Li}_2\text{OH}$  radical was predicted<sup>1</sup> to be thermodynamically stable. The experimental studies on dilithium hydroxide are very limited indeed. The first observation of the  $\text{Li}_2\text{OH}$  radical was made by Kudo et al.<sup>2</sup> in a supersonic beam by photoionization time-of-flight mass spectrometry. The adiabatic ionization energy of  $\text{Li}_2\text{OH}$  was then determined experimentally to be 4.053 eV.<sup>3</sup> To our knowledge, no information is available in the literature on spectroscopic properties of the  $\text{Li}_2\text{OH}$  radical and related species. Therefore, the structure and bonding of dilithium hydroxide is known only from theoretical sources. Schleyer<sup>1</sup> predicted the structure and dissociation energies of  $\text{Li}_2\text{OH}$  at the Hartree–Fock (HF) and Møller–Plesset second-order (MP2) levels of theory using the split-valence double- $\zeta$  (3-21G and 6-31G\*) basis sets. The  $\text{Li}_2\text{OH}$  radical was found to have  $C_{2v}$  symmetry, with the LiO bond dissociation energy of about 40 kcal/mol. Kudo et al.<sup>2–4</sup> investigated the structure and energetics of various  $\text{Li}_n(\text{OH})_{n-1}$  ( $n = 2–5$ ) clusters by the density functional theory (DFT), with the hybrid functional B3LYP, and by the MP2 approach using the split-valence triple- $\zeta$  basis set 6-311+G\*\*. The ionization energies were determined at the B3LYP, MP2, and coupled-cluster CCSD(T) levels of theory. In particular, the adiabatic ionization energy of  $\text{Li}_2\text{OH}$  was predicted<sup>4</sup> to be 4.35, 4.03, and 4.07 eV by the B3LYP, MP2, and CCSD(T) methods, respectively. The structure and vibrational harmonic frequencies of the  $\text{Li}_2\text{OH}^+$  cation were determined by Solomonik and Pogrebnaya<sup>5</sup> at the HF/DZP level. The  $\text{Li}_2\text{OH}^+$  and  $\text{Li}_2\text{OH}^-$  ions were investigated by Nemukhin and Stepanov<sup>6</sup> at the HF and configuration–interaction CISD levels using the 6-31G basis set.

The present study was undertaken with the aim of predicting the structure and spectroscopic properties of dilithium hydroxide by the ab initio approach and of assisting future experimental studies. We present a detailed and accurate characterization of the potential energy surfaces and vibrational–rotational energy levels of dilithium hydroxide and its cation. The molecular

parameters are determined using highly correlated wave functions, calculated with large correlation-consistent spd fgh basis sets, and taking into account core–electron correlation effects. We hope that the results reported here may serve as a useful guide for detecting dilithium hydroxide and related species by vibrational–rotational spectroscopy and for assigning their spectra.

## II. Method of Calculation

The molecular parameters of the  $\text{Li}_2\text{OH}$  radical were calculated using the partially spin restricted open-shell coupled-cluster method including single and double excitations and a perturbational correction due to connected triple excitations, RCCSD(T).<sup>7–9</sup> The coupled-cluster wave function was based on high-spin restricted Hartree–Fock (RHF) molecular orbitals as a reference function. Calculations for the  $\text{Li}_2\text{OH}^+$  cation were performed using the closed-shell coupled-cluster method, CCSD(T).<sup>10,11</sup> The one-particle basis sets employed were the correlation-consistent polarized valence basis sets, cc-pVnZ.<sup>12–14</sup> The quality of the basis sets ranged from triple- to quintuple- $\zeta$  ( $n = \text{T, Q, 5}$ ). The largest basis set, cc-pV5Z, consists of a (14s8p4d3f2g1h)/[6s5p4d3f2g1h] set for lithium and oxygen and a (8s4p3d2f1g)/[5s4p3d2f1g] set for hydrogen. Only the spherical harmonic components of the d through h polarization functions were used. In the valence correlation treatment, the 1s-like core orbitals of the lithium and oxygen atoms were excluded from the active space. The core-related correlation effects were investigated using the correlation-consistent polarized core-valence basis set of triple- and quadruple- $\zeta$  quality, cc-pCVTZ, and cc-pCVQZ.<sup>14,15</sup> These basis sets are obtained by augmenting the corresponding standard cc-pVnZ basis sets with sets of tight functions. The largest basis set, cc-pCVQZ, is augmented with a (3s3p2d1f) set for lithium and oxygen, and it consists thus of a (15s9p5d3f1g)/[8s7p5d3f1g] set. The calculations were performed using the MOLPRO–2006 package of ab initio programs.<sup>16</sup>

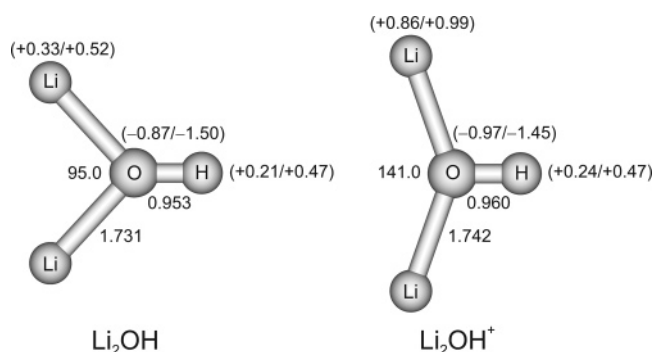
Vibrational harmonic force fields of  $\text{Li}_2\text{OH}$  and  $\text{Li}_2\text{OH}^+$  were calculated using the coupled-cluster methods, RCCSD(T) and CCSD(T), respectively. The anharmonic force fields of these species were determined using the spin unrestricted and restricted Møller–Plesset second-order methods, UMP2 and MP2, respectively.<sup>17</sup> These calculations were performed using the Gaussian 03 package of ab initio programs.<sup>18</sup> The spectro-

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**TABLE 1.** Equilibrium Molecular Parameters of the X<sup>2</sup>A<sub>1</sub> Li<sub>2</sub>OH Radical and the X<sup>1</sup>A<sub>1</sub> Li<sub>2</sub>OH<sup>+</sup> Cation, Determined Using the CCSD(T) Method and Various cc-pVnZ Basis Sets

	cc-pVTZ	cc-pVQZ	cc-pV5Z
X <sup>2</sup> A <sub>1</sub> Li <sub>2</sub> OH			
<i>r</i> (LiO) (Å)	1.7495	1.7460	1.7461
<i>r</i> (OH) (Å)	0.9562	0.9544	0.9542
∠(LiOH) (deg)	132.29	132.25	132.31
Li⋯Li <sup>a</sup> (Å)	2.5884	2.5848	2.5825
energy + 90 (hartree)	-0.721593	-0.748979	-0.757437
<i>r</i> (LiO) <sup>b</sup> (Å)	1.7553	1.7479	1.7468
<i>r</i> (OH) <sup>b</sup> (Å)	0.9573	0.9548	0.9544
∠(LiOH) <sup>b</sup> (deg)	132.59	132.37	132.35
Li⋯Li <sup>b</sup> (Å)	2.5846	2.5827	2.5819
energy + 90 <sup>b</sup> (hartree)	-0.727999	-0.751227	-0.758331
X <sup>1</sup> A <sub>1</sub> Li <sub>2</sub> OH <sup>+</sup>			
<i>r</i> (LiO) (Å)	1.7643	1.7583	1.7576
<i>r</i> (OH) (Å)	0.9623	0.9604	0.9602
∠(LiOH) (deg)	108.98	109.22	109.30
Li⋯Li (Å)	3.3368	3.3206	3.3176
energy + 90 (hartree)	-0.573579	-0.600312	-0.608581
<i>r</i> (LiO) <sup>b</sup> (Å)	1.7656	1.7590	1.7579
<i>r</i> (OH) <sup>b</sup> (Å)	0.9632	0.9607	0.9603
∠(LiOH) <sup>b</sup> (deg)	109.04	109.31	109.31
Li⋯Li <sup>b</sup> (Å)	3.3380	3.3201	3.3180
energy + 90 <sup>b</sup> (hartree)	-0.579337	-0.602342	-0.609422

<sup>a</sup> The calculated distance between the Li atoms. <sup>b</sup> With the aug-cc-pVnZ basis set for oxygen.



**Figure 1.** Predicted equilibrium structures of the ground electronic state of Li<sub>2</sub>OH and Li<sub>2</sub>OH<sup>+</sup>. Bond lengths are given in Å and angles in degrees. Atomic charges, in *e*, determined by the Mulliken/NBO population analyses of the HF/cc-pCVQZ wave functions are given in parentheses.

scopic constants were then determined using the second-order perturbational approach<sup>19,20</sup> as implemented in the Gaussian 03 package.

### III. Results and Discussion

As expected from the previous studies,<sup>1–4</sup> the Li<sub>2</sub>OH radical and the Li<sub>2</sub>OH<sup>+</sup> cation were calculated to have a planar, Y-shaped equilibrium structure of C<sub>2v</sub> symmetry. The ground electronic state of Li<sub>2</sub>OH and Li<sub>2</sub>OH<sup>+</sup> was found to be <sup>2</sup>A<sub>1</sub> and <sup>1</sup>A<sub>1</sub>, respectively. The calculated structural parameters and total energies of both molecules are given in Table 1. The LiO bond lengths of Li<sub>2</sub>OH and Li<sub>2</sub>OH<sup>+</sup> are predicted here to be substantially shorter than those determined by Tanaka et al.<sup>4</sup> at the MP2/6-311+G(d,p) level of theory, being 1.758 and 1.770 Å, respectively; the OH bond length and the valence LiOH angle are quite alike. Following the Mulliken and natural-bond-orbital (NBO)<sup>21</sup> population analyses of the HF electron density, the LiO and OH bonds of Li<sub>2</sub>OH and Li<sub>2</sub>OH<sup>+</sup> are strongly polarized (see Figure 1 below), with a large negative charge of about -0.9*e* (Mulliken) and -1.5*e* (NBO) on the oxygen atom. To allow the one-particle basis set to adjust to such a charge

**TABLE 2.** Equilibrium Molecular Parameters of the X<sup>2</sup>A<sub>1</sub> Li<sub>2</sub>OH Radical and the X<sup>1</sup>A<sub>1</sub> Li<sub>2</sub>OH<sup>+</sup> Cation, Determined Using the CCSD(T) Method with the cc-pCVTZ and cc-pCVQZ Basis Sets<sup>a</sup>

	V	A	A - V <sup>b</sup>
X <sup>2</sup> A <sub>1</sub> Li <sub>2</sub> OH			
cc-pCVTZ			
<i>r</i> (LiO) (Å)	1.7526	1.7399	-0.0127
<i>r</i> (OH) (Å)	0.9569	0.9562	-0.0007
∠(LiOH) (deg)	132.55	132.71	0.16
Li⋯Li (Å)	2.5822	2.5569	-0.0253
energy (hartree)	-90.731823	-90.871586	-0.139763
cc-pCVQZ			
<i>r</i> (LiO) (Å)	1.7472	1.7319	-0.0153
<i>r</i> (OH) (Å)	0.9550	0.9542	-0.0008
∠(LiOH) (deg)	132.38	132.50	0.12
Li⋯Li (Å)	2.5813	2.5538	-0.0275
energy (hartree)	-90.752823	-90.902763	-0.149940
X <sup>1</sup> A <sub>1</sub> Li <sub>2</sub> OH <sup>+</sup>			
cc-pCVTZ			
<i>r</i> (LiO) (Å)	1.7629	1.7504	-0.0125
<i>r</i> (OH) (Å)	0.9629	0.9623	-0.0006
∠(LiOH) (deg)	109.03	109.14	0.11
Li⋯Li (Å)	3.3331	3.3073	-0.0258
energy (hartree)	-90.583150	-90.721747	-0.138597
cc-pCVQZ			
<i>r</i> (LiO) (Å)	1.7584	1.7432	-0.0152
<i>r</i> (OH) (Å)	0.9608	0.9601	-0.0007
∠(LiOH) (deg)	109.31	109.47	0.16
Li⋯Li (Å)	3.3190	3.2870	-0.0320
energy (hartree)	-90.603924	-90.752511	-0.148587

<sup>a</sup> With the aug-cc-pCVnZ basis set for oxygen. <sup>b</sup> A difference between the value determined correlating all of the electrons (*A*) and the value determined correlating only the valence electrons (*V*).

distribution, we investigated the effects of diffuse functions. Additional calculations were performed with the augmented correlation-consistent basis sets aug-cc-pVnZ for oxygen,<sup>22</sup> and the results are also given in Table 1. For the largest basis set cc-pV5Z, enlargement of the basis set gives a negligible effect, the largest changes being observed for the LiO bond length and amounting to about 0.0005 Å. Although both lithium and hydrogen atoms have a single valence electron (2*s* and 1*s*, respectively), the LiO bond of Li<sub>2</sub>OH and Li<sub>2</sub>OH<sup>+</sup> is determined to be at equilibrium nearly twice as long as the OH bond.

As shown in Table 1 and Figure 1, ionization of the Li<sub>2</sub>OH radical changes significantly its equilibrium structure. The LiO bonds lengthen by about 0.01 Å and the valence LiOH angle increases from 95° for the neutral species to as much as 141° for the cation. This is consistent with the results of the Mulliken and NBO population analyses, which both indicate that the additional charge is localized on the lithium atoms. Upon ionization, the total charge of each Li atom of Li<sub>2</sub>OH increases by about +0.5*e*, whereas the total charges of the oxygen and hydrogen atoms change by less than 0.1*e*. On the basis of a simple electrostatic model, the Li<sub>2</sub>OH<sup>+</sup> cation can thus be expected to have a nearly T-shaped equilibrium structure. Apparently, the electrostatic interaction—Coulomb repulsion of the lithium atoms plays a dominant role in determining the shape of the Li<sub>2</sub>OH<sup>+</sup> cation.

The effects of core-electron correlation on the molecular parameters of the Li<sub>2</sub>OH radical and the Li<sub>2</sub>OH<sup>+</sup> cation were investigated by comparing the values determined in calculations correlating only the valence electrons (*V*) with those when all of the electrons were correlated (*A*), both calculations performed with the same core-valence basis set. The results obtained with the cc-pCVTZ and cc-pCVQZ basis sets are given in Table 2.

Inclusion of the core-related effects substantially decreases the equilibrium LiO bond length. The  $A - V$  change for the LiO bond is larger by a factor of  $\sim 20$  than that for the OH bond.

The best estimate of the equilibrium structural parameters can be determined by adding the changes due to the core-related correlation  $A - V$ , determined with the cc-pCVQZ basis set, to the values determined with the largest valence basis set cc-pV5Z. For the Li<sub>2</sub>OH radical, the equilibrium structural parameters are estimated in this way to be  $r_e(\text{LiO}) = 1.731 \text{ \AA}$ ,  $r_e(\text{OH}) = 0.953 \text{ \AA}$ , and  $\angle_e(\text{LiOH}) = 132.5^\circ$ . For the Li<sub>2</sub>OH<sup>+</sup> cation, the equilibrium structural parameters are  $r_e(\text{LiO}) = 1.742 \text{ \AA}$ ,  $r_e(\text{OH}) = 0.960 \text{ \AA}$ , and  $\angle_e(\text{LiOH}) = 109.5^\circ$ . It is worth noting that the estimated values are equivalent to those computed with the cc-pCVQZ basis set while correlating all of the electrons. Considering convergence of the computed values with the one-particle basis set size and the effects of approximations inherent to the CCSD(T) approach, we estimate the uncertainty in the calculated bond lengths and angles to be about  $\pm 0.001 \text{ \AA}$  and  $\pm 0.1^\circ$ , respectively. The predicted molecular parameters of the ground electronic state of Li<sub>2</sub>OH and Li<sub>2</sub>OH<sup>+</sup> are summarized in Figure 1. For both molecules, the predicted equilibrium LiO bond length is substantially longer than that found for dilithium oxide, Li<sub>2</sub>O, determined experimentally<sup>23</sup> and theoretically<sup>24</sup> to be 1.611 and 1.616  $\text{ \AA}$ , respectively. For lithium hydroxide, LiOH, the equilibrium LiO bond length was determined experimentally to be 1.5776  $\text{ \AA}$ .<sup>25</sup> The predicted equilibrium OH bond length of Li<sub>2</sub>OH and Li<sub>2</sub>OH<sup>+</sup> is slightly longer than that of 0.949  $\text{ \AA}$  found for LiOH.<sup>25</sup> The distance between the lithium atoms is predicted to be 2.552 and 3.284  $\text{ \AA}$  for Li<sub>2</sub>OH and Li<sub>2</sub>OH<sup>+</sup>, respectively. For the lithium dimer and its cation, Li<sub>2</sub> and Li<sub>2</sub><sup>+</sup>, the bond length is calculated at the CCSD(T)/cc-pCVQZ level to be 2.6758 and 3.1010  $\text{ \AA}$ , respectively, as compared with the experimental values<sup>26,27</sup> of 2.6730 and 3.12  $\text{ \AA}$ . The predicted equilibrium structural parameters and atomic charge distribution suggest that the Li<sub>2</sub>OH radical can be considered as two ionically bound units Li<sub>2</sub><sup>+</sup> and OH<sup>-</sup>. However, because the Li $\cdots$ Li distance of Li<sub>2</sub>OH is determined to be even shorter than that for the Li<sub>2</sub> molecule, there is strong Li–Li bonding enhanced by the OH<sup>-</sup> moiety. This is similar to the bonding nature of dilithium oxide in its first singlet excited state <sup>1</sup>B<sub>1</sub>.<sup>28</sup> The Li<sub>2</sub>O molecule was found both experimentally<sup>28</sup> and theoretically<sup>29</sup> to be strongly bent in the <sup>1</sup>B<sub>1</sub> state, with the valence LiOLi angle of about 100°. On the other hand, the Li<sub>2</sub>OH<sup>+</sup> cation can be considered as the OH<sup>-</sup> group almost embedded in the Li<sub>2</sub><sup>+</sup> metallic “cluster”.

To determine the shape of the potential energy surfaces of the Li<sub>2</sub>OH radical and the Li<sub>2</sub>OH<sup>+</sup> cation, the vibrational harmonic force fields and frequencies were calculated at various CCSD(T)/cc-pVnZ levels of theory. The results are given in Table 3. The 6 normal modes of the dilithium hydroxide molecule include 3 modes of A<sub>1</sub> symmetry ( $\omega_1 - \omega_3$ ), 1 mode of B<sub>1</sub> symmetry ( $\omega_4$ ), and 2 modes of B<sub>2</sub> symmetry ( $\omega_5, \omega_6$ ), classified according to the C<sub>2v</sub> symmetry group. The assignment of the vibrational modes is based on the potential energy distribution (PED). All of the normal vibrations of both species appear to be characteristic. Among the normal modes involving the lithium atoms, the higher energy mode has LiO stretching character, whereas the lower energy mode has LiOH bending character. It is worth noting that the assignment of the B<sub>2</sub> normal modes of the Li<sub>2</sub>OH<sup>+</sup> cation is reversed. The effects of core–electron correlation on the harmonic frequencies are illustrated in Table 4. As could be expected, almost all of the harmonic frequencies of Li<sub>2</sub>OH and Li<sub>2</sub>OH<sup>+</sup> increase modestly upon inclusion of these effects. The changes due to the core-related

**TABLE 3. Vibrational Harmonic Frequencies (in cm<sup>-1</sup>) of the X<sup>2</sup>A<sub>1</sub> Li<sub>2</sub>OH Radical and the X<sup>1</sup>A<sub>1</sub> Li<sub>2</sub>OH<sup>+</sup> Cation, Determined Using the CCSD(T) Method and Various cc-pVnZ Basis Sets<sup>a</sup>**

assignment	cc-pVTZ	cc-pVQZ	cc-pV5Z
X <sup>2</sup> A <sub>1</sub> Li <sub>2</sub> OH			
$\omega_1$ (OH stretch)	3911.6	3925.8	3929.0
$\omega_2$ (sym LiO stretch)	685.4	693.3	694.2
$\omega_3$ (sym LiOH bend)	270.9	270.3	270.3
$\omega_4$ (out-of-plane)	282.3	277.2	277.3
$\omega_5$ (asym LiO stretch)	649.9	659.8	660.7
$\omega_6$ (asym LiOH bend)	510.4	518.4	519.2
X <sup>1</sup> A <sub>1</sub> Li <sub>2</sub> OH <sup>+</sup>			
$\omega_1$ (OH stretch)	3830.9	3847.2	3850.4
$\omega_2$ (sym LiO stretch)	615.2	624.2	625.0
$\omega_3$ (sym LiOH bend)	224.1	225.3	225.5
$\omega_4$ (out-of-plane)	345.8	346.9	347.0
$\omega_5$ (asym LiOH bend)	798.4	802.8	803.6
$\omega_6$ (asym LiO stretch)	726.8	734.6	736.0

<sup>a</sup> With the aug-cc-pVnZ basis set for oxygen.

**TABLE 4. Vibrational Harmonic Frequencies (in cm<sup>-1</sup>) of the X<sup>2</sup>A<sub>1</sub> Li<sub>2</sub>OH Radical and the X<sup>1</sup>A<sub>1</sub> Li<sub>2</sub>OH<sup>+</sup> Cation, Determined Using the CCSD(T) Method with the cc-pCVTZ and cc-pCVQZ Basis Sets<sup>a</sup>**

	V	A	A - V <sup>b</sup>
X <sup>2</sup> A <sub>1</sub> Li <sub>2</sub> OH			
cc-pCVTZ			
$\omega_1$	3904.8	3911.0	6.2
$\omega_2$	684.9	693.5	8.6
$\omega_3$	271.3	275.4	4.1
$\omega_4$	280.9	287.2	6.3
$\omega_5$	650.7	657.8	7.1
$\omega_6$	509.5	507.2	-2.3
cc-pCVQZ			
$\omega_1$	3925.3	3932.0	6.7
$\omega_2$	693.6	704.4	10.8
$\omega_3$	270.4	275.2	4.8
$\omega_4$	277.6	283.8	6.2
$\omega_5$	660.3	669.6	9.3
$\omega_6$	518.4	516.2	-2.2
X <sup>1</sup> A <sub>1</sub> Li <sub>2</sub> OH <sup>+</sup>			
cc-pCVTZ			
$\omega_1$	3824.9	3829.7	4.8
$\omega_2$	615.1	622.2	7.1
$\omega_3$	224.4	226.3	1.9
$\omega_4$	346.5	350.7	4.2
$\omega_5$	799.6	800.4	0.8
$\omega_6$	726.6	734.1	7.5
cc-pCVQZ			
$\omega_1$	3846.8	3851.9	5.1
$\omega_2$	624.5	634.0	9.5
$\omega_3$	225.3	227.6	2.3
$\omega_4$	346.9	350.9	4.0
$\omega_5$	802.3	803.7	1.4
$\omega_6$	735.0	744.8	9.8

<sup>a</sup> With the aug-cc-pCVnZ basis set for oxygen. <sup>b</sup> A difference between the value determined correlating all of the electrons (A) and the value determined correlating only the valence electrons (V).

correlation  $A - V$ , determined with the cc-pCVQZ basis set, were used to correct the harmonic frequencies determined in the valence-only calculations (Table 3). As for the equilibrium structural parameters, the corrected values appeared to be equivalent to those computed with the cc-pCVQZ basis set while correlating all of the electrons.

The anharmonic force fields of the Li<sub>2</sub>OH radical and the Li<sub>2</sub>OH<sup>+</sup> cation were determined using the spin unrestricted and restricted Møller–Plesset second-order methods, respectively, in conjunction with the cc-pCVQZ basis set and correlating all



**TABLE 5. Vibrational Harmonic ( $\omega$ ) and Fundamental ( $\nu$ ) Frequencies (in cm<sup>-1</sup>) of the X<sup>2</sup>A<sub>1</sub> Li<sub>2</sub>OH Radical and the X<sup>1</sup>A<sub>1</sub> Li<sub>2</sub>OH<sup>+</sup> Cation, Determined at the MP2/cc-pCVQZ<sup>a</sup> Level of Theory**

mode	$\omega$	$\nu$	$\nu - \omega$
X <sup>2</sup> A <sub>1</sub> Li <sub>2</sub> OH			
1	3944.0 (0.2) <sup>b</sup>	3763.6	-180.4
2	694.5 (51)	671.7	-22.8
3	279.4 (14)	271.2	-8.2
4	287.9 (106)	288.4	0.5
5	660.3 (44)	634.1	-26.2
6	510.6 (11)	510.6	0.0
X <sup>1</sup> A <sub>1</sub> Li <sub>2</sub> OH <sup>+</sup>			
1	3862.5 (44)	3688.8	-173.7
2	623.8 (57)	608.8	-15.0
3	226.8 (135)	224.6	-2.2
4	350.6 (219)	352.4	1.8
5	801.7 (40)	794.2	-7.5
6	730.8 (328)	709.5	-21.3

<sup>a</sup> With the aug-cc-pCVQZ basis set for oxygen, all of the electrons correlated. <sup>b</sup> Infrared transition intensities are given in parenthesis, in km/mol.

**TABLE 6. Predicted Vibrational Harmonic ( $\omega$ ) and Fundamental ( $\nu$ ) Frequencies (in cm<sup>-1</sup>) of the Isotopomers Li<sub>2</sub>OH, Li<sub>2</sub>OD, and <sup>6</sup>Li<sub>2</sub>OH**

	Li <sub>2</sub> OH	Li <sub>2</sub> OD	<sup>6</sup> Li <sub>2</sub> OH
$\omega_1$	3936	2866	3936
$\omega_2$	705	699	747
$\omega_3$	275	274	293
$\omega_4$	284	220	284
$\omega_5$	670	644	704
$\omega_6$	517	398	521
$\nu_1$	3755	2770	3755
$\nu_2$	682	676	722
$\nu_3$	267	266	283
$\nu_4$	284	220	285
$\nu_5$	644	621	676
$\nu_6$	517	397	521

of the electrons. For X<sup>2</sup>A<sub>1</sub> Li<sub>2</sub>OH, spin contamination of the UHF wave function appeared to be minimal indeed, with the total-spin-square ( $S^2$ ) expectation value of 0.7503. The predicted harmonic ( $\omega$ ) and fundamental ( $\nu$ ) frequencies of both species are given in Table 5. The harmonic frequencies determined with the MP2 method differ from those determined with the CCSD(T) method by about 10 cm<sup>-1</sup>. As shown by differences between the  $\omega$  and  $\nu$  values, the largest anharmonicity is observed for the OH stretching mode. Anharmonicity of the LiO stretching modes is significantly smaller, with the  $\nu - \omega$  differences being by an order of magnitude smaller than that of the OH stretching mode.

Table 6 summarizes the vibrational constants—the harmonic and fundamental frequencies predicted for the three isotopomers of dilithium hydroxide, namely Li<sub>2</sub>OH, Li<sub>2</sub>OD, and <sup>6</sup>Li<sub>2</sub>OH. The predicted anharmonicity constants  $\chi$  of these species are given in Table S1 of the Supporting Information. The harmonic frequencies  $\omega$  were calculated at the CCSD(T)/cc-pV5Z level of theory and corrected for the core–electron correlation effects determined at the CCSD(T)/cc-pCVQZ level. The anharmonicity corrections  $\nu - \omega$ , determined at the MP2/cc-pCVQZ level, were then added yielding the fundamental frequencies  $\nu$ . For the main isotopic species Li<sub>2</sub>OH, the harmonic and anharmonic zero-point energy (ZPE) is calculated to be 3191 and 3151 cm<sup>-1</sup>, respectively. The vibrational constants predicted for the cations Li<sub>2</sub>OH<sup>+</sup>, Li<sub>2</sub>OD<sup>+</sup>, and <sup>6</sup>Li<sub>2</sub>OH<sup>+</sup> are given in Tables 7 and S2. For the main isotopic species Li<sub>2</sub>OH<sup>+</sup>, the harmonic and anharmonic ZPE is calculated to be 3306 and 3271 cm<sup>-1</sup>, respectively. Considering convergence of the computed values

**TABLE 7. Predicted Vibrational Harmonic ( $\omega$ ) and Fundamental ( $\nu$ ) Frequencies (in cm<sup>-1</sup>) of the Isotopomers Li<sub>2</sub>OH<sup>+</sup>, Li<sub>2</sub>OD<sup>+</sup>, and <sup>6</sup>Li<sub>2</sub>OH<sup>+</sup>**

	Li <sub>2</sub> OH <sup>+</sup>	Li <sub>2</sub> OD <sup>+</sup>	<sup>6</sup> Li <sub>2</sub> OH <sup>+</sup>
$\omega_1$	3856	2806	3856
$\omega_2$	635	633	680
$\omega_3$	228	226	240
$\omega_4$	351	293	355
$\omega_5$	805	746	812
$\omega_6$	746	588	776
$\nu_1$	3682	2714	3682
$\nu_2$	620	618	663
$\nu_3$	226	224	238
$\nu_4$	353	294	357
$\nu_5$	798	724	802
$\nu_6$	725	583	755

**TABLE 8. Predicted Equilibrium ( $B_e$ ) and Ground-State Vibrationally Averaged ( $B_0$ ) Rotational Constants, and Quartic Centrifugal Distortion Constants (in MHz) of the Isotopomers Li<sub>2</sub>OH, Li<sub>2</sub>OD, and <sup>6</sup>Li<sub>2</sub>OH**

	Li <sub>2</sub> OH	Li <sub>2</sub> OD	<sup>6</sup> Li <sub>2</sub> OH
$A_e$	40 708	34 935	44 139
$B_e$	22 113	22 113	25 792
$C_e$	14 329	13 541	16 279
$A_0$	40 731	35 046	44 185
$B_0$	21 932	21 924	25 565
$C_0$	14 166	13 407	16 085
$\Delta_J$	0.22	0.22	0.30
$\Delta_K$	2.64	1.79	3.16
$\Delta_{JK}$	-0.99	-0.85	-1.27
$\delta_J$	0.09	0.10	0.13
$\delta_K$	0.11	0.06	0.10

with the one-particle basis set size and the effects of approximations inherent to the CCSD(T) and MP2 approaches, we estimate the uncertainty in the calculated frequencies to be about  $\pm 5$  cm<sup>-1</sup>.

The predicted rotational and centrifugal distortion constants of the radicals Li<sub>2</sub>OH, Li<sub>2</sub>OD, and <sup>6</sup>Li<sub>2</sub>OH are given in Table 8. The equilibrium rotational constants were calculated with the best estimates of the equilibrium structural parameters discussed above. The quartic centrifugal distortion constants were calculated for the  $A$ -reduction of the vibration–rotation Hamiltonian using the harmonic force field of dilithium hydroxide determined at the CCSD(T)/cc-pCVQZ level of theory. For the three isotopic species, the anharmonic vibration–rotation interaction constants  $\alpha$  were determined at the MP2/cc-pCVQZ level, and their values are given in Table S3. The constants  $\alpha$  were then used to calculate the ground-state vibrationally averaged rotational constants. The predicted rotational and centrifugal distortion constants of the cations Li<sub>2</sub>OH<sup>+</sup>, Li<sub>2</sub>OD<sup>+</sup>, and <sup>6</sup>Li<sub>2</sub>OH<sup>+</sup> are given in Table 9. It is worth noting that due to large structural changes upon ionization, the Li<sub>2</sub>OH radical and the Li<sub>2</sub>OH<sup>+</sup> cation are distinctly different asymmetric tops. The Li<sub>2</sub>OH<sup>+</sup> cation is a nearly prolate symmetric top, with the asymmetry parameter  $\kappa = -0.98$  (close to the limiting case of  $\kappa = -1$ ). The Li<sub>2</sub>OH radical is strongly asymmetric, with the asymmetry parameter  $\kappa = -0.41$ . Both species have thus qualitatively different patterns of the rotational energy levels and consequently qualitatively different rotational spectra. Uncertainties in the rotational constants can be estimated by varying the predicted equilibrium OH and LiO bond lengths and the LiOH angle within their estimated error bars of  $\pm 0.001$  Å and  $\pm 0.1^\circ$ , respectively. For the main isotopic species Li<sub>2</sub>OH, varying the OH bond length results in a small change of the constant  $A$  of  $\mp 10$  MHz, whereas it does not affect the constants  $B$  and  $C$ . Upon varying the LiO bond length, the constants  $A$ ,  $B$ , and  $C$  change by  $\mp 40$ ,  $\mp 25$ , and  $\mp 15$  MHz,

**TABLE 9. Predicted Equilibrium ( $B_e$ ) and Ground-state Vibrationally Averaged ( $B_0$ ) Rotational Constants, and Quartic Centrifugal Distortion Constants (in MHz) of the Isotopomers  $\text{Li}_2\text{OH}^+$ ,  $\text{Li}_2\text{OD}^+$ , and  $^6\text{Li}_2\text{OH}^+$** 

	$\text{Li}_2\text{OH}^+$	$\text{Li}_2\text{OD}^+$	$^6\text{Li}_2\text{OH}^+$
$A_e$	126 130	93 743	134 945
$B_e$	13 357	13 357	15 579
$C_e$	12 078	11 691	13 967
$A_0$	129 032	95 310	138 102
$B_0$	13 284	13 288	15 490
$C_0$	11 977	11 598	13 843
$\Delta_J$	0.05	0.04	0.06
$\Delta_K$	109.36	41.21	120.48
$\Delta_{JK}$	-2.33	-1.26	-2.73
$\delta_J$	0.01	0.01	0.01
$\delta_K$	0.26	0.23	0.33

respectively. Upon varying the LiOH angle, the constants  $A$ ,  $B$ , and  $C$  change by  $\mp 140$ ,  $\pm 70$ , and  $\pm 10$  MHz, respectively.

To estimate intensities of the rotational transitions, the electric dipole moment of the  $\text{Li}_2\text{OH}$  radical was calculated by the finite-field approach at the RCCSD(T)/cc-pCVQZ level with all of the electrons correlated. For the equilibrium configuration, the total dipole moment of  $\text{Li}_2\text{OH}$  was found to be 2.34 D.

Finally, to assess quality of the calculated electronic structure of dilithium hydroxide, the adiabatic ionization energy of the  $\text{Li}_2\text{OH}$  radical was determined. The total energy difference between the equilibrium configurations of  $\text{Li}_2\text{OH}$  and  $\text{Li}_2\text{OH}^+$  is calculated with the CCSD(T) method to be 4.0453, 4.0514, and 4.0520 eV for the cc-pVTZ, cc-pVQZ, and cc-pV5Z basis sets, respectively. Inclusion of the core–electron correlation effects increases the total energy difference by 0.0317 and 0.0368 eV, as calculated with the CCSD(T) method in conjunction with the cc-pCVTZ and cc-pCVQZ basis set, respectively. Inclusion of the vibrational zero-point energies of both species increases further the total energy difference by 0.0143 and 0.0149 eV within the harmonic and anharmonic approximation, respectively. The best theoretical estimate of the adiabatic ionization energy of the  $\text{Li}_2\text{OH}$  radical is thus obtained to be 4.104 eV, as compared with the experimental value of 4.053 eV.<sup>3</sup>

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**Supporting Information Available:** The vibrational anharmonicity constants  $x$  and vibration–rotation interaction constants  $\alpha$  for the three isotopic species  $\text{Li}_2\text{OH}$ ,  $\text{Li}_2\text{OD}$ , and  $^6\text{Li}_2\text{OH}$ , and for the corresponding cations are given in Tables S1–S4. This material is available free of charge via the Internet at <http://pubs.acs.org>.

## References and Notes

- Schleyer, P. v. R. In *New Horizons of Quantum Chemistry*; Löwdin, P.-O., Pullman, B., Eds.; Reidel: Dordrecht, The Netherlands, 1983.
- Kudo, H.; Hashimoto, M.; Tanaka, H.; Yokoyama, K. *J. Mass Spectrom. Soc. Jpn.* **1999**, *47*, 2.
- Tanaka, H.; Yokoyama, K.; Kudo, H. *J. Chem. Phys.* **2001**, *114*, 152.
- Tanaka, H.; Yokoyama, K.; Kudo, H. *J. Chem. Phys.* **2000**, *113*, 1821.
- Solomonik, V. G.; Pogrebnaya, T. P. *Zh. Struct. Khim.* **1983**, *24*, 39.
- Nemukhin, A. V.; Stepanov, N. F. *Zh. Struct. Khim.* **1987**, *28*, 25.
- Watts, J. D.; Gauss, J.; Bartlett, R. J. *J. Chem. Phys.* **1993**, *98*, 8718.
- Knowles, P. J.; Hampel, C.; Werner, H.-J. *J. Chem. Phys.* **1994**, *99*, 5219.
- Knowles, P. J.; Hampel, C.; Werner, H.-J. *J. Chem. Phys.* **2000**, *112*, 3106.
- Raghavachari, K.; Trucks, G. W.; Pople, J. A.; Head-Gordon, M. *Chem. Phys. Lett.* **1989**, *157*, 479.
- Scuseria, G. E.; Lee, T. J. *J. Chem. Phys.* **1990**, *93*, 5851.
- Dunning, T. H., Jr. *J. Chem. Phys.* **1989**, *90*, 1007.
- Feller, D. Unpublished work.
- Woon, D. E.; Peterson, K. A.; Dunning, T. H., Jr. Unpublished work.
- Woon, D. E.; Dunning, T. H., Jr. *J. Chem. Phys.* **1995**, *103*, 4572.
- MOLPRO, version 2006.1, a package of ab initio programs, Werner, H.-J.; Knowles, P. J.; Lindh, R.; Manby, F. R.; Schutz, M.; Celani, P.; Korona, T.; Rauhut, G.; Amos, R. D.; Bernhardsson, A.; Berning, A.; Cooper, D. L.; Deegan, M. J. O.; Dobbyn, A. J.; Eckert, F.; Hampel, C.; Hetzer, G.; Lloyd, A. W.; McNicholas, S. J.; Meyer, W.; Mura, M. E.; Nicklass, A.; Palmieri, P.; Pitzer, R.; Schumann, U.; Stoll, H.; Stone, A. J.; Tarroni, R.; Thorsteinsson, T.
- Møller, C.; Plesset, M. S. *Phys. Rev.* **1934**, *46*, 618.
- Frisch, M. J.; Trucks, G. W.; Schlegel, H. B.; Scuseria, G. E.; Robb, M. A.; Cheeseman, J. R.; Montgomery, J. A., Jr.; Vreven, T.; Kudin, K. N.; Burant, J. C.; Millam, J. M.; Iyengar, S. S.; Tomasi, J.; Barone, V.; Mennucci, B.; Cossi, M.; Scalmani, G.; Rega, N.; Petersson, G. A.; Nakatsuji, H.; Hada, M.; Ehara, M.; Toyota, K.; Fukuda, R.; Hasegawa, J.; Ishida, M.; Nakajima, T.; Honda, Y.; Kitao, O.; Nakai, H.; Klene, M.; Li, X.; Knox, J. E.; Hratchian, H. P.; Cross, J. B.; Adamo, C.; Jaramillo, J.; Gomperts, R.; Stratmann, R. E.; Yazyev, O.; Austin, A. J.; Cammi, R.; Pomelli, C.; Ochterski, J. W.; Ayala, P. Y.; Morokuma, K.; Voth, G. A.; Salvador, P.; Dannenberg, J. J.; Zakrzewski, V. G.; Dapprich, S.; Daniels, A. D.; Strain, M. C.; Farkas, O.; Malick, D. K.; Rabuck, A. D.; Raghavachari, K.; Foresman, J. B.; Ortiz, J. V.; Cui, Q.; Baboul, A. G.; Clifford, S.; Cioslowski, J.; Stefanov, B. B.; Liu, G.; Liashenko, A.; Piskorz, P.; Komaromi, I.; Martin, R. L.; Fox, D. J.; Keith, T.; Al-Laham, M. A.; Peng, C. Y.; Nanayakkara, A.; Challacombe, M.; Gill, P. M. V.; Johnson, B.; Chen, W.; Wong, M. W.; Gonzales, C.; Pople, J. A. Gaussian 03, revision D.01. Gaussian, Inc.: Pittsburgh, PA, 2003.
- Papoušek, D.; Aliev, M. R. *Molecular Vibrational Rotational Spectra*; Academia: Prague, 1982.
- Gaw, J. F.; Willetts, A.; Green, W. H.; Handy, N. C. In *Advances in Molecular Vibrations and Collision Dynamics*; Bowman, J. M., Ed.; JAI Press: Greenwich, CT, 1990.
- Carpenter, J. E.; Weinhold, F. *J. Mol. Struct. (THEOCHEM)* **1988**, *169*, 41.
- Kendall, R. A.; Dunning, T. H., Jr.; Harrison, R. J. *J. Chem. Phys.* **1992**, *96*, 6796.
- Bellert, D.; Winn, D. K.; Breckenridge, W. H. *J. Chem. Phys.* **2002**, *117*, 3139.
- Koput, J.; Peterson, K. A. *J. Chem. Phys.* **2002**, *116*, 9255.
- Higgins, K. J.; Freund, S. M.; Klempner, W.; Apponi, A. J.; Ziurys, L. M. *J. Chem. Phys.* **2004**, *121*, 11715.
- Coxon, J. A.; Melville, T. C. *J. Mol. Spectrosc.* **2006**, *235*, 235.
- Bernheim, R. A.; Gold, L. P.; Tipton, T. *J. Chem. Phys.* **1983**, *78*, 3635.
- Bellert, D.; Winn, D. K.; Breckenridge, W. H. *J. Chem. Phys.* **2003**, *119*, 10169.
- Koput, J. Unpublished work.